DISSOCIATIVE SHOCKS IN THE NEIGHBORHOOD OF ORION IRc2 TRACED WITH ATOMIC CARBON

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ABSTRACT

We present Caltech Submillimeter Observatory 10" resolution maps of the ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ line of neutral atomic carbon (809.3435 GHz) and the J = 7-6 line of CO (806.6518 GHz) covering a region of 2' × 4'.5 around Orion IRc2 in the BN/KL Nebula. The two lines were observed with the same receiver simultaneously, and therefore the two maps have zero relative pointing error. The atomic carbon (C I) emission does not peak toward the IRc2 position and displays a very different spatial distribution with respect to most molecular species. Moderately strong red wings are detected in C I around the region defined by the high-velocity CO gas with a spatial distribution that delineates a shell of ~20" radius. We propose that CO dissociation in the J-type shocks produced by the interaction of the IRc2 outflows with the surrounding quiescent gas could play a significant role in this enhancement of atomic carbon in that shell.

Subject headings: ISM: individual (Orion IRc2) — ISM: molecules — line: profiles — radio lines: ISM — submillimeter

1. INTRODUCTION

Atomic carbon is an important coolant of the neutral interstellar medium (Keene 1995), plays a key role in interstellar chemistry, and is an important tool for probing intermediate regions between diffuse interstellar matter and dense molecular clouds. The first detection of its 492 GHz ${}^{3}P_{1} \rightarrow {}^{3}P_{0}$ fine-structure transition was obtained by the Kuiper Airborne Observatory (KAO; Phillips et al. 1980). The ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ transition (809 GHz) was first detected with the University of Hawaii 88 inch (2 m) telescope (Jaffe et al. 1985). Both lines were first discovered in the Orion molecular cloud. Data from the Cosmic Background Explorer (Wright et al. 1991) have shown that C I emission is widely seen throughout our Galaxy and is generally bright in the Galactic plane. The C I/CO abundance ratio has been confirmed to be high (0.1-1) in nearby galaxies (Stutzki et al. 1997; Gérin & Phillips 2000). Its detailed distribution in individual Galactic clouds has become a subject of study in the past years, and the rapid gain in submillimeter receiver sensitivity has allowed great progress from both ground-based and satellite platforms.

In the chemically rich and complex environment of the Orion massive-star formation region, C I is expected to be a good choice to trace the dissociation of the abundant CO molecule and its possible association with shocks. Phillips & Huggins (1981) had noticed the trend for atomic carbon not to be as strongly peaked at cloud cores compared with molecules and particularly CO. White & Padman (1991) pointed out important differences between the spatial distributions of both species in the Orion region and, in particular, that C I lies mostly in a shell-like structure around IRc2. The first CO J = 7-6 (the highest frequency line of CO that is regularly observed from ground-based observatories) map around Orion IRc2 was obtained with the KAO at 98" resolution (Schmid-Burgk et al. 1989). Recent observations with ~ 10 m dishes have improved the angular resolution to $\sim 15''$ (Wilson et al. 2001). Large areas of the integral-shaped region of Orion have been recently mapped in the two lowest fine-structure lines of C I (Plume et al. 2000; Yamamoto et al. 2001; Ikeda et al. 2002), but with modest angular resolutions: 1'.4–2'.2. The region around IRc2 was mapped at the James Clerk Maxwell Telescope using the ${}^{3}P_{1} \rightarrow {}^{3}P_{0}$ line (492 GHz) with an angular resolution of 14" by White & Sandell (1995).

In this work we present the spatial distribution of atomic carbon and carbon monoxide in a region of $2' \times 4'5$ centered on Orion IRc2 using the Caltech Submillimeter Observatory (CSO). Our choice for tracing CO is through its J = 7-6 line originating at an energy level of 156 K, which should therefore be especially well adapted for tracing the moderate- and high-velocity "plateau" structure. A C I transition $({}^{3}P_{2} \rightarrow {}^{3}P_{1})$ is very close in frequency (809 GHz), so these two lines are both accessible with similar atmospheric transmission and can be mapped simultaneously at virtually the same resolution. The observations and results are presented in § 2. The data are discussed in § 3.

2. OBSERVATIONS AND RESULTS

The observations were performed with the 10.4 m telescope of the CSO at the summit of Mauna Kea (Hawaii) in 2000 September. The receiver was a helium-cooled SIS mixer operating in double-sideband mode. For details about the receiver, see Kooi et al. (2000). Two acousto-optic spectrometers with 0.5 and 1.0 GHz coverage providing velocity resolutions of 0.174 and 0.55 km s⁻¹, respectively, were used simultaneously during the observations. The pointing was checked every hour using Saturn and was determined to be good within 3", for a half-power beamwidth of ~10" at 806 GHz. As both lines were observed with the same receiver tuning, relative pointing problems are eliminated.

Atmospheric precipitable water vapor was in the range of 0.7–1.0 mm. The system temperature was ~4000 K. The observations were performed using the on-the-fly procedure with the antenna moving at an angular speed of $1'' \text{ s}^{-1}$ across a row made of 11 boxes of $10'' \times 10''$ each. A calibration scan and an off-position scan 120 s long, a degree away from IRc2 in

azimuth, were performed at the end of each row. A total of 297 positions were observed with this procedure.

The calibration is achieved by assuming that the effective temperature of the atmosphere equals the ambient temperature. The data have been corrected for the difference between these temperatures in order to obtain accurate antenna temperatures, T_A^* (see Pardo et al. 2005). At 806 GHz the correction factor is as high as 15% at the elevation of the source for $\tau_{\rm atm,\,zenith} \sim 1.5$.

Figure 1 shows CO J = 7-6 and C I ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ spectra at selected positions. Our CO 7–6 data have much better signalto-noise ratios than previous observations obtained at similar angular resolutions due to system temperatures being in our case at least a factor of 4 lower and similar integration times. Toward the position of IRc2 we can clearly see that there is emission between about -90 and 80 km s⁻¹. The shape of our signal toward this position coincides quite well with the CO J = 4-3 observation of Wilson et al. (2001). The velocity extent of the signal quickly decreases as the distance to IRc2 increases; 1' away from IRc2 the line is already less than 30 km s⁻¹ wide. Only the extended ridge is then detected. CO and C I emission also displays line wings at position $\Delta \alpha = -10^{"}$, $\Delta \delta = -80^{"}$, i.e., in the direction of the Orion cross-outflow (Schmid-Burgk et al. 1990).

Our C I ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ data, in addition to having much better angular resolution, also significantly improve the signal-tonoise ratio and the velocity resolution, with respect to those presented in Yamamoto et al. (2001). The Plume et al. (2000) and Ikeda et al. (2002) data of the lower ${}^{3}P_{1} \rightarrow {}^{3}P_{0}$ transition have a poor angular resolution, $\approx 2'$ in both cases, and comparison with our results is difficult. The White & Sandell (1995) map of the lower transition is the only one useful for comparison with our C I map (see next section).

3. DISCUSSION

The features that have been revealed in the Orion Nebula by the large number of millimeter and submillimeter observations available are the following: First, a dense and warm clump called the "hot core." It is quite compact (10"), with temperatures estimated in the range 400-500 K and densities reaching 10^7 cm⁻³. However, the detection of CO transitions with upper rotational levels up to J = 43 reveal the presence of a much hotter component with T = 1500-2000 K and CO column densities on the order of 2×10^{17} cm⁻² (Sempere et al. 2000). Second, a moderate velocity "plateau" consisting of gas heated to temperatures ranging from 80 to 200 K distributed in a region covering $\sim 40''$ centered on IRc2 with an average density below 3×10^6 cm⁻³. Lines arising from this component typically have half-power line widths of about 40–50 km s⁻¹ (see, for example, HCN and HDO in Pardo et al. 2001), although high-velocity gas is seen in CO in the velocity ranges -60 to -40 and 50 to 80 km s⁻¹. Finally, a much extended structure called the "ridge" with densities of $\sim 10^4 - 10^5$ cm⁻³, velocity widths of $\sim 8 \text{ km s}^{-1}$, and temperatures still reaching 50-60 K due to heating by the UV field from the Trapezium stars. The emission from and the radiative transfer through all these components together give a quite complex output.

The high-velocity (40–60 km s⁻¹) CO 7–6 emission peaks on IRc2 and also on the H₂ peak 2 (see Sugai et al. 1994 as a reference for Orion's molecular hydrogen peaks). The extremely high velocity CO gas (\pm 60–80 km s⁻¹) is found along a filament connecting the positions of IRc2 and H₂ peak 1. This filament is not necessarily a jet, but the result of heating



FIG. 1.—CO J = 7-6 spectra at selected positions around Orion IRc2 ($\alpha = 5^{h}35^{m}14^{s}5$, $\delta = -5^{\circ}22'31''$ [J2000.0]). The line at -105 km s⁻¹ is the ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ transition of atomic carbon from the image sideband. Therefore, the velocity scale for it has to be recentered and reversed.

by the embedded sources of the gas inside the cavity eroded by Orion's outflows. Although the filament is mainly seen at extreme velocities, emission at all velocities exists along it. This spatial distribution of the high-excitation CO emission is rather different from the one traced by its low-J lines. Moreover, the spatial distribution of the extreme negative and positive velocities in the CO J = 7-6 emission looks very similar to that of water vapor at 183.3 GHz, as observed by Cernicharo et al. (1994). Taking into account that this line of water is extremely sensitive to the kinetic temperature, its spatial coincidence with the CO 7-6 line reinforces the idea that the latter is tracing a warm-gas component, probably related to the one traced by the high-excitation lines of CO (up to J = 43) observed with the Infrared Space Observatory by Sempere et al. (2000) and the rovibrational lines of CO and H₂O observed by González-Alfonso et al. (2002).

The spatial distribution of C I obtained in this work provides a clearer picture than previously. Toward IRc2 the line peaks at ~9–10 km s⁻¹ and has a width at half-intensity on the order of 5 km s⁻¹. However, at other positions (for example, see position $\Delta \alpha = 30''$, $\Delta \delta = 0''$ in Fig. 1) moderate wings are clearly detected and an interesting velocity structure is revealed (see Fig. 2). The integrated line intensity delineates a doublepeaked structure with the two peaks defining a line quite close to that of the central section of the integral-shaped large-scale structure of the Orion Nebula. The highest C I ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ intensity



FIG. 2.—*Top two rows*: Color map at the top left shows the observed total integrated line intensity of the ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ transition of atomic carbon. The other panels of the top two rows show the CO J = 7-6 integrated intensity at different velocity ranges (indicated at the top of each panel). The integrated intensity for some maps has been divided by a scale factor (indicated in the panels as necessary) to maintain the same color scale for all maps (the units for all color maps in this figure are K km s⁻¹). The panels corresponding to velocities from 0 to 20 km s⁻¹ also show as red contours the integrated intensity in the corresponding interval of atomic carbon (first contour and step are both 2 K km s⁻¹). The positions of BN, IRc2, and H₂ peaks 1 and 2 (see text) are marked for reference and identified in the 60–80 km s⁻¹ map. *Bottom row*: The left panel is a zoom of the integrated intensity map of atomic carbon over the central region of Orion IRc2. The contours correspond to the blue (from -40 to -20 km s⁻¹) high-velocity CO J = 7-6 emission. Note that the atomic carbon intensity has its maximum at the edge of the region where the high-velocity CO gas is present. The right panels show the atomic carbon intensity at different velocities (indicated in each panel) over a range of 0.5 km s⁻¹. The contours correspond to the emission of the J = 12-11 transition of HC₃N as observed with the 30 m IRAM telescope and kindly provided by A. Rodríguez-Franco (see Rodríguez-Franco et al. 1992; first contour and step are both 0.5 K km s⁻¹).

is found virtually in coincidence with the shocked molecular hydrogen peak 1. The southwest C I peak is mainly blueshifted emission, whereas the northeast, more intense peak is redshifted with respect to the peak velocity of CO, ~9 km s⁻¹. The ${}^{3}P_{1} \rightarrow {}^{3}P_{0}$ (492 GHz) maps presented in White & Sandell (1995) also have two peaks, northeast and southwest of IRc2, although slightly more separated than ours. The most prominent peak in their case is also the northeast one, which has the shape of an arc, relatively similar to what we see. The velocity-binned maps of ${}^{3}P_{1} \rightarrow {}^{3}P_{0}$ suggest that C I delineates a shell-like structure around IRc2, a result that is now much better supported by our velocity-binned maps of ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ presented in Figure 2. The C I contours for the 15–20 km s⁻¹ velocity range still seem to provide a hint of Orion's outflow.

In our case, especially for the northeast C I peak between 10 and 12 km s⁻¹ but also for the southwest one, a puzzle-like complementarity with the high-velocity (both wings) CO J = 7-6 emission suggests that shocks created by the interaction of the central outflow with the extended gas component dissociate CO and therefore cause the enhancement of atomic carbon in the shock fronts. The radius of the delineated C I shell is ~20". Combined with this, in the large-scale distribution of atomic carbon other factors should play a role, such as

intrinsic differences in column density and temperature and the effect of UV photons from the nearby ionizing O and B stars.

In order to check the effect that volume density differences may have on the observed C I spatial distribution, it is helpful to compare with a high-density tracer. The HC₃N J = 12-11data from Rodríguez-Franco et al. (1992) have been kindly provided by the authors. The comparison of C I and HC₃N J = 12-11 spatial distributions at different velocities is also shown on Figure 2. For the velocity range 7–9 km s⁻¹, the HC₃N emission is well centered on IRc2, but when the northeast C I peak appears at v > 9 km s⁻¹, the HC₃N peak moves to be almost coincident with it, at least in the range 9.5–11.0 km s⁻¹. This fact indicates that the stronger intensity of the northeast C I peak with respect to the southwest one is probably due to a larger gas column density there.

The column density of atomic carbon can be estimated from its two fine-structure lines. The data from White & Sandell (1995) for the ${}^{3}P_{1} \rightarrow {}^{3}P_{0}$ C I line can be compared with ours for ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ with a similar spatial resolution. As discussed by these authors, the excitation temperatures of the atomic carbon lines are not very sensitive to the physical parameters of the different gas components of the Orion cloud. One could assume that the atomic carbon lines are thermalized with an excitation tem-

perature close to that of the carbon monoxide lines. The observed intensities for both lines suggest excitation temperatures close to 100 K. We obtain column densities for C I in the range $(0.2-2) \times 10^{18} \text{ cm}^{-2}$ using the equation proposed by Keene et al. (1985), but correcting for the line strength ratio between the ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ and ${}^{3}P_{1} \rightarrow {}^{3}P_{0}$ lines. The deduced abundance ratio between atomic carbon and carbon monoxide varies therefore between ~0.01 and 0.35 according to the results found for CO by White & Sandell (1995) using C¹⁸O and ¹³CO, and our own calculations using an LVG code with $T_{\kappa} = 100$ K (CO column density in the range 6 \times 10¹⁸ to 3 \times 10¹⁹). It is reasonable to think that the maximum is located at the shell of enhanced C I emission and that the minimum is located in the direction of IRc2, although new detailed observations of optically thin lines of isotopic CO at similar resolutions would be necessary to confirm this point.

The fact that C I is the result of the dissociation, driven by photons or shocks, of a molecule with a large abundance in molecular clouds (CO) and having lines in the submillimeterwave region, where appropriate resolutions can be achieved with available telescopes, has provided this unique opportunity to trace this dissociation in the vicinity of Orion IRc2. Other atomic species (i.e., O I with SOFIA and *Herschel*) and radicals

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should be mapped at similar or higher resolutions in order to check whether or not the effect of dissociative shocks in the neighborhood of Orion IRc2 is the main thing responsible for what has been presented in this letter.

4. SUMMARY

CO 7–6 and C I ${}^{3}P_{2} \rightarrow {}^{3}P_{1}$ mapping around Orion IRc2 with high angular resolution (~10") has been achieved with no relative pointing problems. The data reveal a double-peaked structure for the distribution of C I in the neighborhood of Orion IRc2, indicating a shell of radius ~20", that shows a striking complementarity with the distribution of high-velocity CO 7–6. It is proposed that this structure could be evidence of a shockdriven dissociation of CO produced as the molecular outflow impinges on the surrounding gas. Local differences in the general column density have been explored using a high-density tracer (HC₃N), which seems to explain why one of the C I peaks (to the northeast) is stronger.

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